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ATMOSPHERIC TRACE CONSTITUENTS MEASUREMENTS
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ANALYSIS OF SOLAR SPECTRA IN THE MIDDLE ULTRAVIOLET AND VISIBLE FOR ATMOSPHERIC TRACE CONSTITUENTS MEASUREMENTS

NASA Langley Research Center

Contract NSG 1405

Final Report

15 March 1977 - 15 May 1980

Submitted by

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20 May 1980



Detailed description of the previous work under this project has been included in the annual and semi-annual reports, for the period 15 March 1977 - 15 May 1979. During that period, analysis was made of selected ultraviolet and visible spectra from the data obtained during two balloon flights made in 1977. A compilation of the best available spectral absorption coefficients data for O₃, NO₂, HNO₂, NO₃ and ClO has been completed, and used for generating synthetic spectra in the 2800 - 7000 Å region for comparisons with the flight data. With the synthetic spectra atmospheric NO₂ features were identified on the sunset spectra and used for derivation of NO₂ mixing ratio altitude profiles from both flights. The results were described in references 1-3. The present report covers the final period of 15 May 1979 - 15 May 1980.

Data from the 10/10/79 UV balloon flight has been analyzed in an attempt to determine the amount of atmospheric OH. The flight yielded solar spectra at ~0.05Å resolution in the 3060-3090Å region. Numerous good scans were obtained during ascent and from float altitude (~33 km) during sunset.

The balloon data has been investigated for possible features of atmospheric OH superimposed on the solar OH⁽⁴⁾ and atomic lines. The expected atmospheric absorptions are quite small, (5-7) so the search was conducted by ratioing high sun scans to low sun scans, similar to our work on the NO₂ profiles. (2) The low sun scans, however, are

strongly affected by the O₃ attentuation. Optimal filtering ⁽⁸⁾ of the data prior to the ratioing is essential for this work and has been implemented.

Typical spectra obtained during the 10/10/79 flight were presented in the last semiannual report (9). The resolution of the data (~0.05Å) is close to that obtained in the Göttingen atlas, (10) which is lower than that obtained in the Kitt Peak atlas. (11) The wavelength calibration was done for soan 80 and is based on the Kitt Peak atlas. Wavelength shifts between consecutive scans have been eliminated by cross-correlation programs that were developed for the 1977 UV-visible flights.

Point-by-point ratios of low sun scans to high scans have been formed in the 3070-3085 Å region. It is in this region that the strongest features of the $A^2\Sigma - X^2\Pi(0,0)$ band are expected to occur. These are shown in Figure 1, where OH lines positions (in Å) and intensities (in atm⁻¹cm⁻² at 240K) are marked. The line positions and intensities have been generated by Goldman and Gillis. (12) Figure 1 and Tables 8, 9 of the paper show the OH line parameters at 240K and 4600K. The complete manuscript is given in the Appendix. The largest atmospheric path achieved before total attenuation by O_3 occurred was with zenith angles of ~87°, corresponding to ~0.1 airmass. Examination of these ratioed spectra has failed to reveal any features which can unambiguously be assigned to atmospheric OH. Assuming that we can detect as little as 2% absorption in the ratioed spectrum, an upper limit can

be placed on the vertical column abundance for OH from 33 km altitude of $\sim 6 \times 10^{12}$ molecules/cm². This is within an order of magnitude of the abundances measured by Anderson (5) and Burnett (6, 7).

The estimate of OH abundance from our data is done by considering the weak line approximation, where the equivalent width AW of a line (of a line group) of intensity S is given by

$$\Delta W(cm^{-1}) = S(cm^{-2}atm^{-1})N(atm cm),$$

where N is the absorber amount.

The strongest feature on our spectrum at the atmospheric temperature of 240K is the partially resolved line group

Transition	<u>\(\lambda\)</u>	S(cm ⁻² atm ⁻¹)
Q ₁ (3, 5)	3081.5479	1621
$Q_{P_{21}}^{(3.5)}$	3081.6259	345
P ₁ (1.5)	3081.6677	2060
	Total Inten	sity $S = \overline{4026}$

Assuming we can detect 2% absorption and that our resolution is $0.05\text{\AA}\approx0.05\,\text{cm}^{-1}$, Δ W = (0.02) x (0.5 cm^{-1}) = 0.01 cm⁻¹, then the minimum detectable absorber amount is

$$N = \frac{0.01}{4026} = 2.5 \times 10^{-6} \text{ atm cm}$$

= 2.5 x
$$10^{-6}$$
 (atm cm) x 2.69 x 10^{19} (molec/cm³) x $\frac{273}{240} \approx 8 \times 10^{13}$ molec/cm².

The airmass for the scans used (114-118) is ~ 0.1 and at 33 km, the vertical airmass is $\sim 8 \times 10^{-3}$, so that the corresponding vertical column is $\sim 6 \times 10^{-12}$. It is estimated that higher spectral resolution (say, ~ 0.01 Å) and a more optimal atmospheric path to minimize O_3 absorptions with possibly larger airmass will allow the identification and quantification of OH on such spectra.

Acknowledgment is made to the National Center for Atmospheric Research which is sponsored by the National Science Foundation, for computer time used in this research. Part of the computer programming and analysis was done by Darwin Rolens and James Gillis.

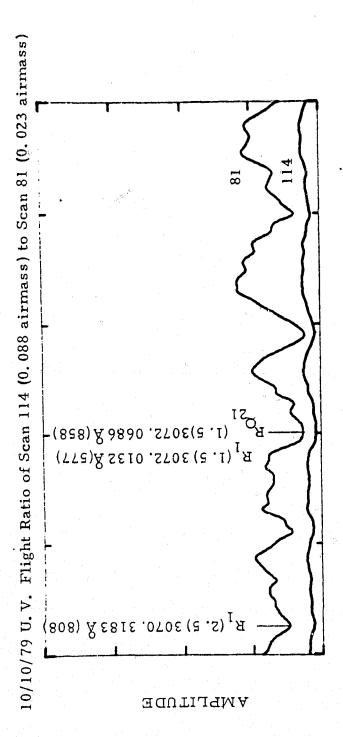
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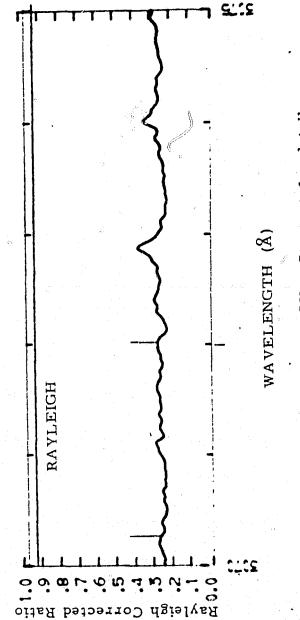
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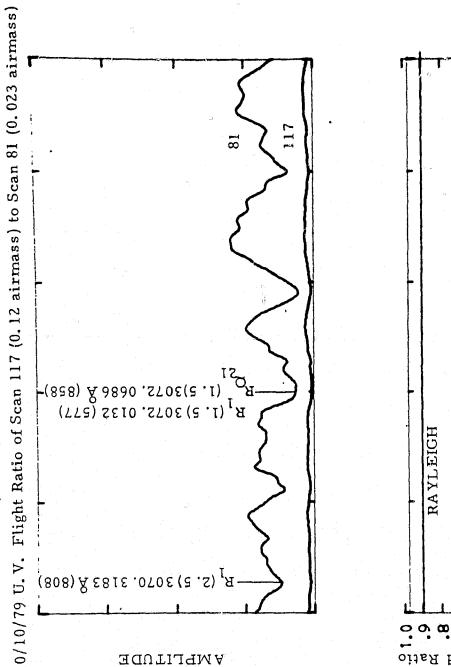
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g. 1. The search for OH. See text for details.



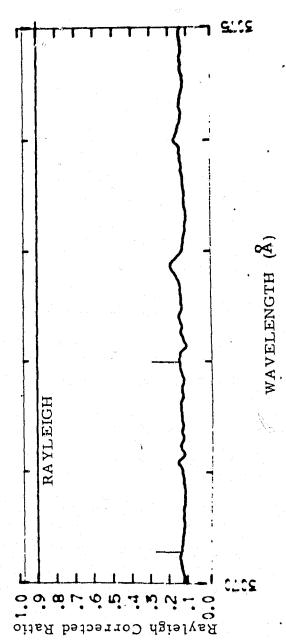
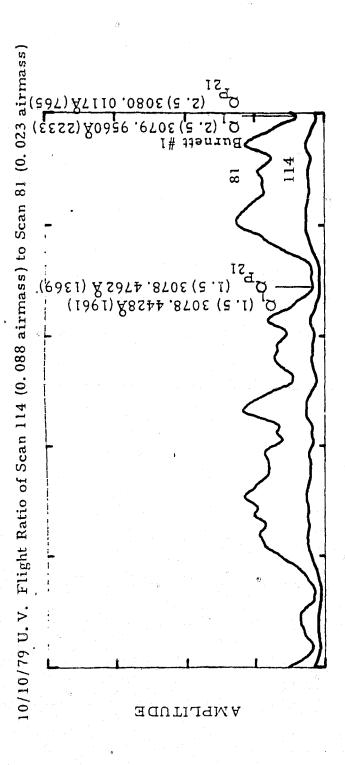
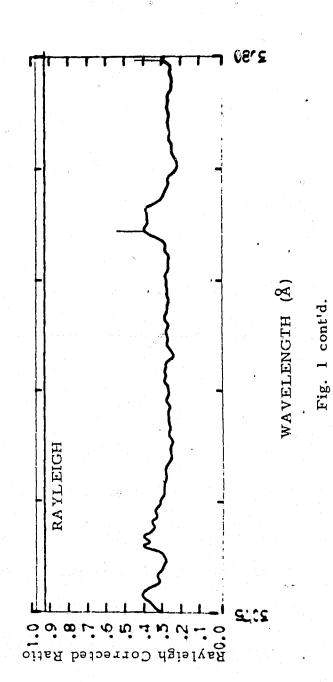


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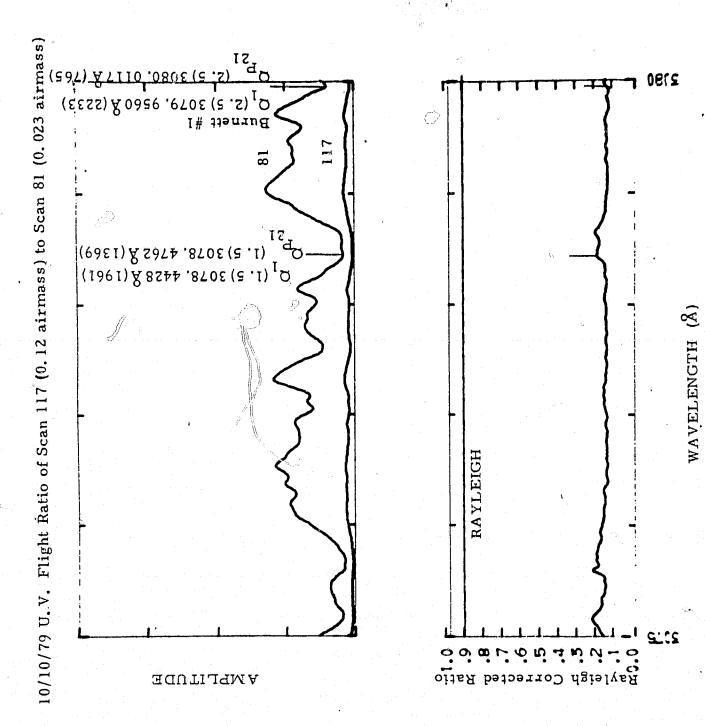
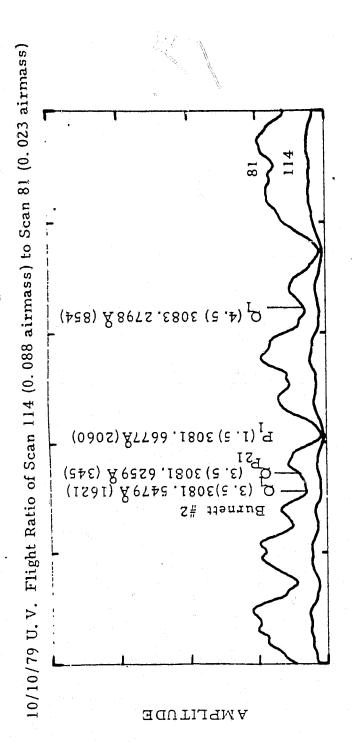
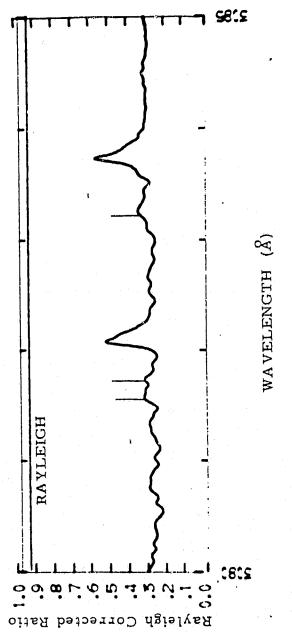
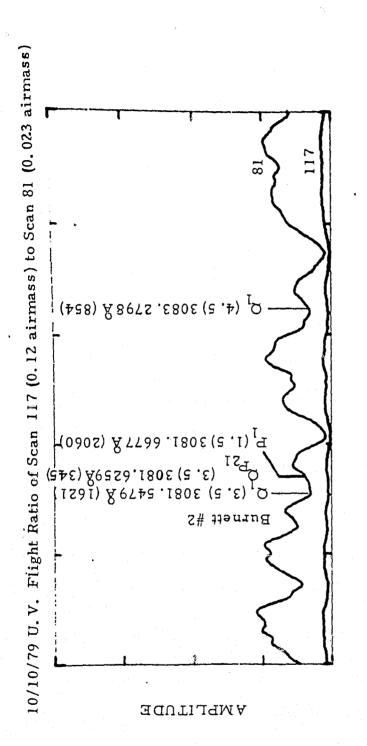


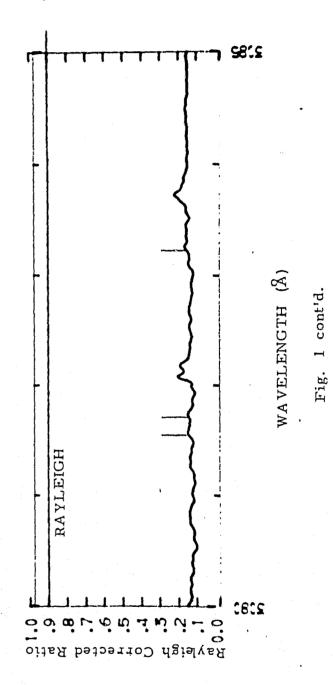
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APPENDIX

SPECTRAL LINE PARAMETERS FOR THE $A^2\Sigma - X^2\pi(0,0)$ BAND OF OH FOR ATMOSPHERIC AND HIGH TEMPERATURES

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Abstract - Individual spectral line parameters including line positions, strengths, and intensities have been generated for the $\Lambda^2\Sigma-X^2\Pi(0,0)$ band of OH, applicable to atmospheric and high temperatures. Energy levels and transition frequencies are calculated by numerically diagonalizing the Hamiltonian. Line strengths are calculated using the dipole matrix and eigenvectors derived from energy matrix diagonalization. The line strengths are compared to those calculated from previously published algebraic line strength formulas. Tables of line parameters are presented for 240 K and 4600 K.

1. INTRODUCTION

The $A^2\Sigma - X^2\Sigma(0,0)$ band of OH in the 3085 Å region has been of interest to quantitative spectroscopists for many years because of its high absorption and emission intensity and convenient wavelength location for spectroscopic probes. The hydroxyl radical is a common by-product of most combustion processes, is present in atmospheric, solar and stellar spectra, and in recent years has been also recognized as an important trace constituent in atmospheric chemistry. Accurate determination of the amount of OH present during spectroscopic experiments depends on precise knowledge of line positions and intensities. Several analyses of spectral line positions for this band are available; among the more important of them are those of Dieke and Crosswhite, who provided the first extensive analysis of the OH UV spectrum, and Destombes et al., 2 who performed elaborate analysis of modern microwave, IR and UV OH data. Intensity (relative and absolute) studies of this band have been reviewed recently by Chidsey and Crosley, 3 who also performed extensive RKR calculations of rotational transition probabilities for the A-X system of OH.

The purpose of this work is to combine the best presently available data and theory to derive accurate quantitative line parameters for the $A^2\Sigma - X^2\Pi(0,0)$ band, applicable to atmospheric and high temperatures. The results are displayed in line parameter tables, and include improved values for the line strength, calculated in intermediate coupling from the energy matrix eigenvectors.

2. LINE PARAMETERS DERIVATION

The OH molecule has an unpaired electron with total electronic angular momentum L = 1 and spin S = $^1/_2$. In the electronic ground state the projection of L along the internuclear axis is Λ = 1. The projection of S along the internuclear axis is $\Sigma = \pm^1/_2$, with a total electronic angular momentum projection $\Omega = \Lambda + \Sigma$. Here, Λ , Σ , Ω are considered as signed quantities, as in the notation of Hougen. The electronic ground state is an inverted $^2\Pi$ state with the $^2\Pi_{1/2}$ ($\Omega = \pm^1/_2$, F2) levels at higher energy than the $^2\Pi_{3/2}$ levels ($\Omega = \pm^3/_2$, F1). The rotational levels for this state are intermediate between Hund's cases (a) and (b). In the $^2\Sigma$ upper state, which is Hund's case (b), $\Lambda = 0$ and $\Omega = \pm^1/_2$ with $J = N\pm^1/_2$. The $^2\Pi$ and $^2\Sigma$ states perturb one another resulting in Λ doubling for each N ($^2\Sigma_{1/2}$) or J ($^2\Pi_{1/2}$, $^3/_2$).

We use the unique perturber approximation described by Destombes et al., 2 to calculate energy levels. This process is restricted to a single vibrational level v in the A $^2\Sigma$ - X $^2\Pi$ subspace. The total angular momentum number F is a good quantum number. For a given F, the J and J+1 levels are weakly coupled by magnetic hyperfine interactions. These interactions are negligible in the calculation of electronic spectra and J may be considered a good quantum number. This gives a 6 x 6 Hamiltonian matrix, the elements of which are given in Table 1. The matrix elements are written in Humd's case (a), with the wave functions represented by $|\Lambda S\Sigma > |J\Omega > = |\Lambda S\Sigma; J\Omega >$, so that

$${}^{2}\Sigma_{\pm^{1}/2}^{+}: |\Lambda S \Sigma; J \Omega \rangle = |0^{+}|^{1}/_{2}\pm^{1}/_{2}; J \pm^{1}/_{2} \rangle$$

$${}^{2}\Pi_{\pm^{1}/2}: |\Lambda S \Sigma; J \Omega \rangle = |\pm^{1}|^{1}/_{2}\mp^{1}/_{2}; J \pm^{1}/_{2} \rangle$$

$${}^{2}\Pi_{\pm^{3}/2}: |\Lambda S \Sigma; J \Omega \rangle = |\pm^{1}|^{1}/_{2}\pm^{1}/_{2}; J \pm^{3}/_{2} \rangle$$

$$(1)$$

The Hamiltonian constants used 5 are shown in Table 2 (these constants give a better fit of the observed spectrum than those of Ref. 2, which do not have a sufficient number of digits retained). The 6 x 6 Hamiltonian matrix may be reduced to two 3 x 3 blocks by the Kronig transformation

$$|J\Omega\delta\rangle = \frac{1}{2} (|\Lambda S\Sigma\rangle |J\Omega\rangle + \delta |-\Lambda S-\Sigma\rangle |J-\Omega\rangle) , \qquad (2)$$

where δ equals s = symmetric = + or a = antisymmetric = -. In this new basis, one of the 3 x 3 blocks contains only matrix elements of the type $\langle J'\Omega's | H | J\Omega s \rangle$ and the other contains only matrix elements of the type $\langle J'\Omega's | H | J\Omega a \rangle$.

After numerical diagonalization of a (3 x 3) Kronig transformed block, the electronic state of each eigenvalue (energy level or, more properly, term value in cm⁻¹) may be determined by noting that the largest eigenvalue belongs to the $^2\Sigma_{1/2}$ state, the intermediate eigenvalue to the $^2\Pi_{1/2}$ state and the smallest eigenvalue to the $^2\Pi_{3/2}$ state. The remaining quantum numbers and parities may be determined from Table 3.

The selection rules for $A^2\Sigma - X^2\Pi(0,0)$ electric dipole allowed transitions are $\Delta J = 0$, ± 1 , $\Delta N = 0$, ± 1 , ± 2 . Parity selection rules require $+ \leftrightarrow -$, $+ \nleftrightarrow +$, and $- \nleftrightarrow -$, which give wavefunction Kronig symmetry selection rules of s \leftrightarrow s and a \leftrightarrow a for $\Delta J = \pm 1$ (P and R branches) and s \leftrightarrow a for $\Delta J = 0$ (Q branches). Transitions are designated by $^{\Delta N}\Delta J_{F^1F^{11}}(J^{11})$ where 'refers to the upper state and "refers to the lower state. These selection rules permit 12 branches, of which 6 are main branches ${^{(P}P_{11}, {^{Q}Q_{11}, {^{R}R_{11}, {^{P}P_{22}, {^{Q}Q_{22}, {^{R}R_{22}})}}}$ and 6 are (weaker) satellite branches ${^{Q}P_{21}, {^{R}Q_{21}, {^{R}R_{21}, {^{P}Q_{12}, {^{Q}R_{12}, {^{Q}P_{12}})}}}$.

The line intensities $S_{i,u}(T)$ (cm⁻¹/atm cm) at temperature T (K) are calculated from 6

$$S_{ku}(T) = \frac{1}{8\pi cv^2} \left(\frac{N}{p}\right) \frac{e^{-1.\Delta_{358E''/T}}}{c_{-3}} A_{v''J''}^{v'J'} (2J'+1) (1-e^{-1.4388v/T}), \qquad (3)$$

where v (=E'-E') is the transition frequency in cm⁻¹, c = 2.99792458 × 10^{10} cm/sec, N is the transition of OH molecules/cm⁵, p is the pressure in atm, $A_{v''J''}^{v''J''}$ is the Einstein A coefficient in sec⁻¹, E'' is the lower state energy (= E₁), and Q_{vR} is the vibration rotation partition function.

Line intensities may be converted from ${\rm cm}^{-1}/{\rm atm}$ cm at T to ${\rm cm/molecule}$ at T by

$$S_{\ell u}(T) \text{ (cm/molecule)} = 3.721963 \text{ x } 10^{-20} \frac{T (K)}{273.16 (K)} S_{\ell u}(T) \text{ (cm}^{-1}/\text{cm atm}).$$
 (4)

The $S_{g,u}(T)$ in cm/molecule are at the population temperature.

We assume ${\bf Q_{VR}}$ may be given by ${\bf Q_{V}Q_{R}}$ where the vibrational partition function ${\bf Q_{V}}$ in the harmonic oscillator approximation is

$$Q_{v} = \frac{1}{1 - e^{-1.4385} - e/T}$$
, (5)

and ω_e is the vibrational harmonic oscillator frequency in cm⁻¹. Huber and Herzberg give $\omega_e = 3737.75$ cm⁻¹. A comparison of the values of Q_v calculated by Eq. (5 and by direct summation of $e^{-E_v/T}$ shows a difference of less than 0.2 percent at 4600 K. The rotational partition

function $Q_{\mathbb{R}}$ is calculated from the actual energy levels as

$$Q_{R} = \sum_{J'',\Omega''} (2J''+1) = -1.-358\Xi''(J''\Omega'')/T \qquad .$$
 (6)

Chidsey and Crosley give the Einstein A coefficient as

$$A_{v''J''}^{v'J''} = \frac{64\pi^4}{3h} p_{v''J''}^{v'J'} S_{J'J''} v^3/(2J'+1) sec^{-1} , \qquad (7)$$

where $p_{V''J''}^{V'J''}$ is the electronic radial transition probability and $S_{J'J''}$ is the rotational line strength. Chidsey and Crosley tabulate relative values of $A_{V''J''}^{V'J''}$ through N'' = 32 for the $A^2\Sigma - X^2\Pi(0,0)$ band in their Table 4. They state that they calculated the line strengths $S_{J'J''}$ based on Earls' formulas with a J dependent spin-orbit coupling parameter A and the rotational constants of Dieke and Crosswhite. We have found that use of Earls' formulas can lead to significant errors at high J in the satellite bands (more about this in the following section). Therefore, we have calculated $S_{J',J''}$ following the method described by Hougen (we describe this in some detail below). Chidsey and Crosley have kindly provided a table of relative $p_{V'',J''}^{V',J'}$ through J = 35.5 prior to publication.

The $\mathbb{A}_{V^{1'}J^{1'}}^{v^{1'}J^{1'}}$ may be put on an absolute basis by noting that the lifetime of a state is

$$\tau_{\mathbf{v}'\mathbf{J}'} = \begin{pmatrix} \Sigma & A_{\mathbf{v}'\mathbf{J}''} \\ \mathbf{v}''\mathbf{J}'' & A_{\mathbf{v}''\mathbf{J}''} \end{pmatrix}^{-1} \text{ sec} . \tag{8}$$

Because Chidsey and Crosley 3 give $A_1^{\circ}/A_0^{\circ}=0.0040$ (here the notation is $A_{v''}^{v'}$), we assume that for the v'=0 vibrational state all vibrational states other than v''=0 make negligible contributions to $\tau_{v',T'}$. The

best available lifetime for the rotationless (N' = 0) v' = 0 state is probably that measured by German, 10 $^{7}_{c}$, $^{1}_{/2}$ = (0.688±0.007)x $^{10^{-6}}$ sec. There are three transitions from the v'' = 0 state to the rotationless v' = 0 state, namely $^{P}_{11}$ (1.5), $^{O}_{12}$ (1.5), and $^{P}_{Q_{12}}$ (0.5). We calculate relative $A_{v''J''}^{v'J'}$ for all J' and J'' of interest by ignoring all constant factors in Eq. (7) and normalize them using Eq. (8).

We form the $S_{J^1J^{11}}$ following the method given by Hougen. ⁴ In addition to the selection rules $\Delta J=0,\pm 1$ and $+\leftrightarrow-$ parity, electric dipole selection rules on Ω give nonzero matrix elements only for $<\Omega\pm 1$ $|\nu_{\mathbf{x}}\pm i\nu_{\mathbf{y}}|\Omega>$ for $\Delta\Omega=\pm 1$ and $<\Omega$ $|\nu_{\mathbf{z}}|\Omega>$ for $\Delta\Omega=0$ where $\nu_{\mathbf{x}}$, $\nu_{\mathbf{y}}$, and $\nu_{\mathbf{z}}$ are electric dipole moment components in the molecule fixed axis system. In the laboratory fixed coordinate system

$$\mu_{Z} = \frac{1}{2} (\alpha_{Zx}^{-i\alpha_{Zy}}) (\mu_{x}^{+i\mu_{y}}) + \frac{1}{2} (\alpha_{Zx}^{+i\alpha_{Zy}}) (\mu_{x}^{-i\mu_{y}}) + \alpha_{Zz}^{-i\alpha_{z}}, \qquad (9)$$

where α_{Zx} , α_{Zy} , and α_{Zz} are the direction cosines between the molecule-fixed and laboratory-fixed coordinate systems. The direction cosine matrix elements are given in Table 4. The μ_x , μ_y , and μ_z are taken to be experimentally determined parameters. For lack of better information, we take $\frac{1}{2}|\mu_x + i\mu_y| = \frac{1}{2}|\mu_x - i\mu_y| = |\mu_z| = 1$, i.e., the electronic perpendicular and parallel transition moments are equal (see below).

Because we assume the electronic and rotational parts of the wavefunction are separable, the electric dipole matrix elements may be written as

$$<\Lambda'S'\Sigma'; J'\Omega'; +_{Z_1} \Lambda S\Sigma; J\Omega>$$

$$= \frac{1}{2} <\Lambda'S'\Sigma' | \mu_{X} + i\mu_{Y} | \Lambda S\Sigma> < J'\Omega' | \alpha_{ZX} - i\alpha_{ZY} | J\Omega>$$

$$+ \frac{1}{2} <\Lambda'S'\Sigma' | \mu_{X} - i\mu_{Y} | \Lambda S\Sigma> < J'\Omega' | \alpha_{ZX} + i\alpha_{ZY} | J\Omega>$$

$$+ <\Lambda'S'\Sigma' | \mu_{Z} | \Lambda S\Sigma> < J'\Omega' | \alpha_{ZZ} | J\Omega> , \qquad (10)$$

where "has been dropped on the lower state quantities. Only one of the three terms in the right hand side of Eq. (10) is nonzero for any allowed transition matrix element. The electric dipole transition matrix is formed in the same basis as was the Hamiltonian matrix, that is, the Hund's case (a) basis. Selection rules in this basis are $\Delta S = 0$, $\Delta \Sigma = 0$, $\Delta \Lambda = 0, \pm 1$, and $\Delta J = 0, \pm 1$. Because $\Delta \Sigma = 0$ and $\Omega = \Lambda + \Sigma$, the selection rule $\Delta \Omega = \Delta \Lambda$ is obtained. As we are interested in $\Sigma = \Pi$ transitions, we set all matrix elements of the type $\langle \Sigma | \mu_Z |^2 \Sigma \rangle$ and $\langle \Pi | \mu_Z |^2 \Pi \rangle$ to zero. The relative phases of the matrix elements are determined by following the prescription of Hougen and Whiting and Nicholls. Following the suggestions of Whiting and Nicholls, we have normalized the line strengths so that

$$\sum_{\Sigma, J'} {}^{S}J'J'' = 2(2S + 1)(2J'' + 1) = 4(2J'' + 1) , \qquad (11)$$

Table 4 reflects this choice of normalization.

The line strength S_J, J'' is formed in intermediate coupling by taking $|< u|\mu_Z| \ell >|^2$ where |u> represents the eigenvector of the $|\Sigma|$ state, |E> represents the eigenvector of the |E| state and $|\mu_Z|$ now represents the 6 x 6 transition matrix. However, the eigenvectors formed during the diagonalization of the energy matrices are in the Kronig transformed basis; therefore, the dipole matrix must also be transformed into this basis. Using the Kronig transformed wavefunctions given by Eq.(1), the Kronig transformed dipole matrix elements have the form

$$$$

$$= \frac{1}{2} \{ \langle J'\Omega' | \mu_{Z} | J\Omega \rangle + \delta \langle J'\Omega' | \mu_{Z} | J-\Omega \rangle \}$$

$$+ \delta' \langle J'-\Omega' | \mu_{Z} | J\Omega \rangle + \delta \delta' \langle J'-\Omega' | \mu_{Z} | J-\Omega \rangle \} , \qquad (12)$$

and, if we let μ_Z^K be the matrix of the < $J^!\Omega^!\delta^!|\mu_Z|J\Omega\delta$ >,

$$\mathbf{S}_{\mathbf{J}^{1}\mathbf{J}^{11}}^{\mathbf{I}} = \left| \langle \mathbf{J}^{1}\Omega_{\mathbf{S}}^{1} \hat{\mathbf{\varepsilon}}^{1} | \mu_{\mathbf{Z}}^{\mathbf{K}} | \mathbf{J}\Omega\delta \rangle \right|^{2}. \tag{13}$$

To our knowledge the $\frac{2}{2}$ Transition matrices have nor been published elsewhere. We show them for the P,Q, and R branches in Tables 5 through 7. The twelve branches correspond to the four (3x3) blocks in the Kronig basis as follows /

For Q-branch lines, the only non-zero(3x3)blocks are the a-s and s-a, while for R- and P-branch lines, the only nonzero(3x3)blocks are the a-a and s-s. When forming line strengths, the appropriate block must be substituted into Eq. (13) for μ_Z^K . For example, to generate the Q₁ line strength we form

$$S_{J^{i}J^{i}}^{Q_{1}} = \left| \langle \psi^{2} | (z \leftarrow s) | \psi^{s} \rangle \right|^{2}.$$

$$2_{\Sigma_{1}/2}$$

$$2_{\Pi_{3}/2}$$
(15)

3. RESULTS AND DISCUSSION

Line strengths, Einstein A coefficients, intensities, and transition frequencies have been calculated for all branches of the $\stackrel{2}{\text{A}} \stackrel{2}{\text{\Sigma-X}} \Pi(0,0)$ OH spectrum through J = 15.5 at 240 K for atmospheric applications and through J = 40.5 at 4600 K for high temperature applications. These calculated values are shown in Tables 8 and 9, respectively. The total band intensities (by summation of the individual lines) are 2.7948 x_0 10 cm⁻¹/atm cm at 240 K and 8.6863×10^{2} cm⁻¹/atm cm at 4600 K. Line intensities are plotted at these two temperatures in Fig. 1. Some caution must be exercised in using high J data from Table 9. The spectroscopic constants used here were determined from data 2 which included transitions through J = 25.5. Although these constants allow prediction of that data to within 0.1 cm -1 maximum error and a standard deviation of ~0.03 cm -1 (hyperfine structure is neglected here), such accuracy cannot be expected for all lines between J = 25.5 and 40.5. Uncertainties in the calculation of energy levels at these high J cause proportionally smaller uncertainties in the energy eigenvectors (wavefunctions) and in quantities calculated using the eigenvectors (line strengths', Einstein A coefficients, and intensities).

Chidsey and Crosley Plist $p_{V''J''}^{V'J'}$ through J = 35.5. We have extrapolated $p_{V''J''}^{V'J'}$ for J = 36.5 through 40.5. Although the dependence of $p_{V''J''}^{V'J''}$ on J is quite linear for P, Q, and R transition probabilities between J = 25.5 and 35.5, extrapolated $p_{V''J''}^{V'J''}$ used to calculate Einstein A coefficients and intensities at higher J must be used with caution.

Although we list four digits for Einstein A coefficients and intensities in Tables 8 and 9, the absolute uncertainties of these quantities cannot be less than 1 percent, because German's 10 $\tau_{0,\frac{1}{2}}$ has a 1 percent quoted uncertainty and because Chidsey and Crosley's 9 $p_{v''J''}^{v'J'}$ are quoted to three significant

digits. The relative uncertainties are limited by the relative accuracy of the $p_{V^{11}J^{11}}^{V^{1}J^{11}}$ and the line strengths. These relative uncertainties should be less than 0.5 percent for J less than 25.5 and are probably less than one percent for J less than 35.5. It should be noted that the present results are based on the assumption that the magnitudes of the parallel and perpendicular electronic transition moment components are equal. To check this assumption we calculated oscillator strengths $f_{l,l}$ according to Penner's Eq. (2-21) and compared them to the tabulated values of Sutherland and Anderson. We found no differences between the two sets of $f_{l,l}$ significant enough to indicate that the magnitudes of the parallel and perpendicular electronic transition moment components are different.

Our calculated line strengths have been checked for accuracy by comparison with the values calculated using Earls' algebraic formulas (these formulas are equivalent to those of Kovacs 13 for 2 E- 1 transitions). When centrifugal and higher order distortion and E- 1 interactions are ignored, our line strengths are identical to those calculated using Earls' formulas. Ignoring these effects does not significantly alter the line strengths for main branch transitions, but does lead to large errors at high J for the weaker satellite branch line strengths. Earls' formulas (which ignore centrifugal and higher distortion and E- 1 interactions) predict consistently smaller line strengths than those calculated by us (which include centrifugal and higher distortion and E- 1 interactions). In the 0 R₁₂ branch, Earls' line strengths range from 93 percent of our line strength at J'' = 20.5; in the 12 R₂₁ branch they range from 78 percent of our line strength at J'' = 39.5. Earls' formulas show errors intermediate in this range for the other satellite branches.

Bennett's 14 line strength formulas, which include P centrifugal distortion, may be expected to give much more accurate results. Based on our

check of Earls' formulas, the line strengths for the $^{\rm S}R_{21}$ branch should have the largest deviation. Bennett's formula predicts line strengths for this branch which are 1 percent high at J = 1.5, decreasing to 6 percent low at J = 25.5, and increasing to 2 percent high at J = 39.5. Examination of the eigenvectors shows that the $^{\rm P}^6$ term, and, to a lesser extent, the $^{\rm 2}\Sigma^{-}\pi$ mixing, can contribute an effect of a few percent to the satellite bands line strengths. Thus, the dominant effect in the deviations from Earls' formulas is due to the centrifugal distortion, which is relatively large in a light molecule such as OH.

The conclusion to be drawn from these comparisons is that both the algebraic formulas or our method give accurate main branch line strengths at all experimentally observed J values. Earls' formulas lead to significant errors at high J in the satellite bands. Bennett's formulas give acceptable satellite branch line strengths for most work. However, when highest accuracy is required, the line strengths from Table 9 should be used.

When our Einstein A coefficients are normalized to the same relative value as those of Chidsey and Crosley, the two sets of values differ by at most ±2 in the last decimal place. As with the line strengths, these differences become important only at high J in the satellite branches where many of the relative Einstein A coefficients are quoted to only one significant digit by Chidsey and Crosley. An additional advantage of our Einstein A coefficients in Tables 8 and 9 for quantitative spectroscopy is that they are absolute rather than relative values.

Although we have chosen to present OH line parameters for temperatures of 240 K and 4600 K, our computer program can generate A Σ -X $\Pi(0,0)$ band line parameters for any temperature. Table 10 may be used with Table 8 or 9 and Eq.(3) and (5) to convert line intensities from these temperatures to any temperature in the 200 - 6000 K range. Intensities so determined should have the

same accuracy as those in Tables 8 and 9. Band intensities at temperatures other than 240 K and 4600 K may be calculated by summing the individual line intensities at the desired temperature. Simpler approximate procedures which directly convert from a band intensity at one temperature to a band intensity at another temperature such as Eq. (7-126) in Penner 6 give errors of approximately 15 percent when band intensities at 240 K and 4600 K are compared.

Acknowledgments - This research was supported in part by NASA/ Langley contract NSG 1405. We thank J.L. Destombes and C. Marliere-Demuynck for kindly providing molecular constants which predict the line positions better than those previously published, and I.L. Chidsey and D.R. Crosley for kindly providing tables of $p_{v''J''}^{v'J'}$ prior to their scheduled publication. Acknowledgment is made to the National Center for Atmospheric Research, which is sponsored by the National Science Foundation, for computer time used in this research.

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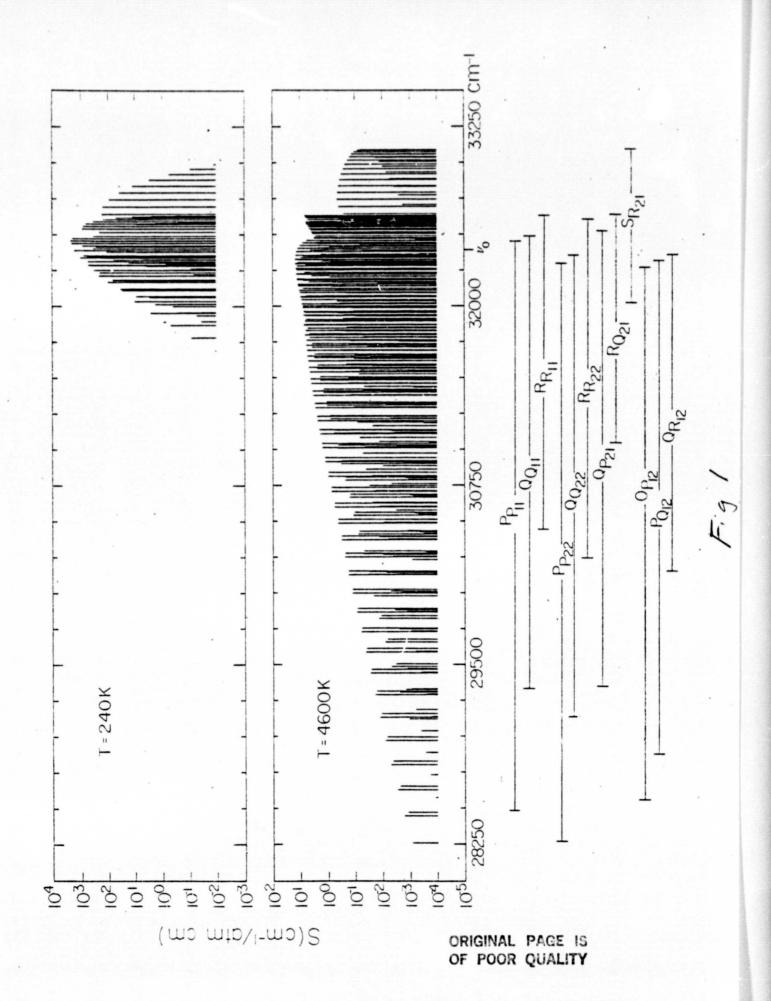
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FIGURE CAPTION

Figure 1. Line intensities and positions for the $A^2\Sigma - X^2\Pi(0,0)$ band of OH. Only lines with intensity greater than 10^{-5} the intensity of the strongest line have been plotted.



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- Table 3. Assignment of quantum numbers to eigenvalues for given J
- Table 4. Nonzero rotational direction cosine matrix elements in $\langle J'\Omega'|\mu_Z|J,\Omega\rangle$ (Hund's case (a))
- Table 5. P-branch ${}^2\Sigma^{-2}\Pi$ transition matrix
- Table 6. Q-branch ${}^2\Sigma^{-2}\Pi$ transition matrix
- Table 7. R-branch ${}^2\Sigma {}^2\Pi$ transition matrix
- Table 8. Line parameters for the $A^2\Sigma X^2\Pi(0,0)$ OH band at 240 K
- Table 9. Line parameters for the $A^2\Sigma X^2\Pi(0,0)$ OH band at 4600 K
- Table 10. Rôtational partition function Q_R vs temperature T

Matrix Element

Value

Notes:

- 1. $x=J+\frac{1}{2}$; $y=[(J-\frac{1}{2})(J+\frac{3}{2})]^{\frac{1}{2}}$
- 2. Matrix elements are unchanged by exchange of initial and final states or by setting Ω to $-\Omega$ in both initial and final states.

Constant	Value * (cm ⁻¹)
B ₂ *	16.9258978
D_{Σ}	$2.0396_{682} \times 10^{-3}$
${\tt H}_{\Sigma}$	97.7 ₆₁₂ x 10 ⁻⁹
γ_{Σ}	$-7.8_{395} \times 10^{-3}$
Vo	32402.056230
A	-139.228 ₃₂₅
ВП	18.5497 ₃₅₄
D _{II}	$1.907_{852} \times 10^{-3}$
H _I	$0.1239_{836} \times 10^{-6}$
A _D	$-0.72_{304} \times 10^{-3}$
<al<sub>+></al<sub>	-151.9226 ₂₁₂
<bl<sub>+></bl<sub>	25.0435440
<dl<sub>+></dl<sub>	$2.6923_{534} \times 10^{-3}$
<a<sub>DL₊></a<sub>	$8.051_{637} \times 10^{-3}$
<hl<sub>+></hl<sub>	$0.166_{605} \times 10^{-6}$

^{*} Values are from reference 5 and are rounded to three figures beyond the standard errors indicated in reference 2.

State		² II·3/ ₂	² 11/ ₂	2Σ 1	1/2
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N	•	$J^{-1}/2$	J+1/2	$J^{-1}/_{2}$	J+1/2
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Antisy	ymmetric block	(-1) ^N	(-1) ^{N+1}	(-1) ^N	

The phases employed here are the same as those used in References 2, 4, and 11.

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$$x = \left[\frac{j-1/2}{2J}, \frac{(j+1/2)}{2J}\right]^{1/2}$$

ii) Kronig basis

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$$v = \left[\frac{(J+^{1}/2)^{2}(2J+1)}{2J(J+1)} \right]^{1/2}$$

ii) Kronig basis

i) Hund's case (a) basis

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$$z = \frac{(J+^{1}/2)(J+^{3}/2)}{2(J+1)}^{1/2}$$

 $y = \left[\frac{(J-1/2)(J+1/2)}{2(J+1)}\right]^{1/2}$

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י ער		ABLAST	2611.	2151.	7.0	539F-E	-811E-2	120E+0	\$1170E+01	6
, ,		86.11	7 (11.	2154	6		.641E-2	-447E+0	. 06416E-01	53
, W		9	-	170°	2	9.696E+01	. 17 1E-	0	.33399E+00	20
ľ,	4.5	01.00	0 1 29 4.	2171.	1	•292E+0	.077E-1	.037E+0	4.83615.00	ij
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r, I	12.5	52.48	P 2 1(12.	2183.	9	.895E-U	2-3202.	375+0	52693E-C1	2
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	5	29.45	, m	2214.	6	.019E	.33 GE-1	1.377E+05	.55810F+00	63
'n		419.0	2 d	2223.	102	.764E-C		.0 40E+0	97759E-01	7.0
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5	11.5	419-98	111.	2226.		.986E-6	•	139E	30295E+6	2 :
ا ا		0.96.56	٠	2225.	102	******************	7.868E-23	•	98938E-0	7
ָּי אַ אָ	יי מינו	289.061	p p 2 2 2.51	7257		3.293E+02			2.14981E+00	16
		64.2	P 1 21 1.	2253.	660	8.761E+01	2.865E-18	Ξ.	10162E-0	16
S	2	9.34	1 2(2.	2253.	660	2.407E+02	•	٠.	577350+0	11
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ı, ı	20.0	ė.	613	261.	3098.8158	3.295E-03	7.007E-22	6.2469E+04	7.41015t-01	
ν π	7.5	0,1	77.	22B1.	960	3.359E+00	, .			12
, IC		7	1 20 7	2283.	960	1.7765-01	5.806E-21	•	976845-0	63
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. 10		078.5	2 21 6	30%	194.62	1.664E+01	•	7	27554E+01	63
5		711.5	9	2306.	94.460	1.106E+00	•	•	.59222E-01	2
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5	20.1	24.8		32323.000	797	6.436E+01	2.105E-10 1.802E-19	6 - 252E + 05	05547708F-01	V P
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ر د د	3.5	201.922	. 5 5 5 6 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	2340.56	92.	3	-487E	.955E+U	00 + 32 40 5 T + 4	
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	INTENSITY CHAMOL COLLE	40.0 K	. 40 3E-1	1956-1	-365E-2	. 54 3E-1	. 284E-1	5-137F-18	. 369E-1	- 80 TE-1	.450E-1	1-3604	1.1916-18	. 897E-1	-39 60	. 793E-	-737E-	5.302F-17	. 39 2E-	.501E-	.301E-	- 44 / 6E -	115E-	.860E-	-375E-	• 84 IE -	. 80 4E -	14.5E-	-019E-	• 196E-	.037E-	5.871E-25	• 63 2E -	.117E-	-3906·	- 32 3E -	. 091E-	. 891E-	8.491E-21	1.0/0E-21	1. 10 2E-17	37.96-		
	INTENSITY PM-2 ATM-6	2 = 1	_ (€	.711596	641E+0	.776E+0	.984E+0	.53/E+U	559E+0	584E+0	*055E+0	•679E+0	6425+0	.721E+0	.250E+0	.541E+0	OFOE PO	1.625F+03	.982E+0	.647E+0	.233E+0	*369E+G	.472F+0	098E+0	.205E-C	.300E+0	.5/5E+U .771F+P	.616E-0	.118E+0	.715E+0 .n76E+0	.396E-0	~	.056E+0	4756+0	.723E-0	.628E+G	337E-6	.190E-0	.596E-0	0-3067•	15E+0	339E-0	54E-0	
re o, cont.	HAVELENGTH .		374	1.27	9.900	3.871	9.855	3089-8156	9.734	7.486	. 342	5.394	3035,3217		-	.	j.	1 2 2	0.00	0.00	670		077	07.4	073	072	2 / 0	071	0.20	070	0.40	3069.6780	690	0.68	068.	969	5000	190	190	99	0.00	990	990	
TABLE	FREQUENCY VAC CM-1 AT	_	349.1	2350.	1324.1	356	356.5	166.4522	2355.8	32379.405	2380.9	2390.0	2012	2415.4	24.22.5	2423.	2440.5	* ·	2455.5	2457.9	5454.5	2474.1	32474553	2517.4	2531	2540.4	2541-9	2551	2553.6	2559	2561.2	2	2572.	2577.5	2578.	2582.0	2 2 2 2 2 2 2 3 2 3 3 3 3	2585.9	2587.3	2588.5	32591.026	2599.7	32603.003	
	TRANSITION		0 2 21 3.	9, 1, 2(3	p 2 11 7.	0 2 21 1.	0 2 2(2.	R 1 2 1.	0 1 10 7	P 2 11 6.	n 1 16 6.	P 1.11 2.	P 2 11 5	R 2 21	P 2 11 4.	n 1 11 4.	P 1 11 1.	P 2 11 3.	R 2 2(1.	P 2 11 2.	0 1 16 2.	p 2 11 1.	0 2 26 2	R 2 2(3.	R 2 2(14.	R 2 2(4.	0 2 1(1.	R 2 2(13.	R 2 21 5.	0 2 11 2.	n 2 1(15.	R 2 2(12	0 2 26 6.	R 1 1 3.	R 2 2(11).	R 2 21 7.	1111 2 0	R 1 1(14.	R 2 2(8.	R 2 2(9.	0 2 11 4	0 2 1(13.	R 1 1113.	٥ <u>,</u>
		•	c	c	C	C .	c (- -	. •	•	د	•	6 6	7 E		~	٠.	c c	· ~	·	c	C (~ •	(: ≥	~	2	∝ 0	. c .	¥ .	~ 0	ν α	~	~ (× œ	~	ez :	~ •	* 12	~	2 1	2 6	× 6×	8	
	LOJER ENERGY	3	9.27	29.27	29.09	87.49	88.76		29.09	12.69	69.21	83.71	•	26.64	5.90	5.90	0.00		87.75	55.×92	3.92	0.5	90.0	29.4	1.36	08-18	000	41.67	324.52	3.71	341.84	316	77.85	01.92	880.52	67.61	010 .E0	810.75	693.29	054.35	55-10	3311.982	11.98	
	36		i.	3.5			•	 			•.		ny n ny n	• i	3		.	-, -		~	۲,	.	 .	4 100	3	÷	-		72		7.5	12	, ٺ	, F	11.	:	-	• • •	6	9		# P7	13.	
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ORGINAL PAGE 18

	MBCB	151	152	153	150	154	156	151	158	159	150	141	152	163	164	<u>ان</u>	166	195	15.3	163	170	171	172	17.3	174	175	176	177	178	179	180	
	LTHE STRF115TII NUMBER	. 25 10 aF + 0	.6132FF	.90527F-3	.17227E+0	1-19353F + 01	. 67187E+0	9	0. • iii	-09040E+0	.72852E+0	0-3696	.01913610	.518176-0	9.06007F+00	.780505.0	.8270FF 0	-87007E-	. 38634F-	- 4893AF-	- 533 19E-	- 596965.	-101214-	-31653F-	-1538RF-	- 99663F -	-85026E-	.71673E-0	• 59626E-	-362	.39187E-0	
	EINSTEIN A L Seg-1	.527E+0	.352E+0	. 975E+D	.9n9E+0	2.697E+05	. 502E+0	17E+0	4.835E+04	.T19E . 0	• 604E + 0	.742E+0	.953E 1 0	. 274E+0	-	. F.71E+0	.767E . 0	. 305E + 0	.756E+0	.569E+0	. 221E+0	7.3E+	• 569E • 0	.319E+0	1.1156+04	. 494E+0	8.1436+03	•035E+0	. 6.110E+03	5-338E+03	.687E+0	
e 	THTENSITY CHZHOLEGULE 40.0 K	1.3446-18	. 90 2E-1	.733E-2	3.259E-19	-2	• 55 SE-	7	?	2 1 E	. 847E-1	.323E-2	9	•	1.4.32E-21		~	. 34.2E-1	5.135E-18	ნ.	. 20 PE-1	3.81AE-19	-49 RE-2		. 014E-2	3. BB 9E-22		7	2.412E-25	.37	•	
	CH-Z ATH-1	.111E+C	-	.365E-0	.965E+0	-	756	.420	.942E+0	346E-	9E+0	.045E-0	.055E-0	.897E-0	4.380E-02	.293E+0	.545E-0	.63%E+0	570	.967E+0	•695E±0	1.60	•905E+0	5.769E-01	.217E-0	.189E-0	0-3	.061E-0	7.377E-06	4.200E-07	1.965E-08	
TABLE 8, cont.	MAVELENGTH R ANGSTPOHS	066.118	165.	0.65.	665.112	3065,0879	190	964.	190	990	190	063	063	69.0	063	063,720	063.	062.	057.	053.	048.57	044.330	0.0	ė	033.431	3030.4901	02	0.25.75	0	022.645	021.73	
TAB	FREDUCHCY VAC CH-1 AT	2665.0	2605.5	2612.9	2615.7	32616,012	2617.4	22.3	2623.6	2625.1	2625.5	2627.9	2628.3	2629.8	2630.5	2633.5	()	2543.2	25.11.5	2744.5	2792.7	2038.3	2001.3	2.020.5	2956.3	2988.3	3016.2	3039.9	3059.2	3.9	3083.	
	TRANSITION	C	R 1 1(5.	0 2 1112.	0 2 16	R R 1 1(12.5)	R 1 11 6.	0 2 1(11.	0 2 11 7.	R. 1 1(11)	2	0 2 1(10.	c	C	R 1 161	R 1 11 B.	R 1 11 9.	R 2 11 1.	e :	R 2 11 3.	R 2 11 4.	œ	R 2 11 6.	≃	R 2 11 8.	R 2 11 9.	~	R 2 1(11.	R 2 1112.	R 2 1(13.	R 2 1(1	
	LOWER ENERGY VAC CH-1	543.575 -			767.456 +	2846.089 +	• 4.5	12	.73	13.61	1:920	2015.036 +	1321,252 +	1650.790 -	2015.036 +	1.25	56 . 79		-92	202.370 +	355.900 -	544.809 +	12.	1029.092 +	1324.291 -	-		-	52.48	3319,355 +	19.15	
	J. 96 U.	ند	-	12.	ڡؙ	.5.12.5		11.		11.		10	.5 6.5		.5 10.5				•	.5 3.5	.5 4.5	.5 5.5	9	.5 7.5	~	6	10.	11.	12.	5 13.	=	

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240.0 K 2.79460E+04 CH-2 ATH-1 AT T = 4.21460E+07 SEC-1 THE INTEGRATEL INTENSITY FOR THE BAND IS THE EINSTEIN A COEF FOR THE BAND IS

9.13939E-16 CM/HOLECULE

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JG 76	LOWER ENERGY VAC CH-1	TRANSITION	FREQUENCY V AC CH-1	MAVELENGTH AIR ANGSTROMS	INTENSITY CH-Z ATH-1	INTENSITY CH/MOLECULE	EINSTEIN A L	LINE STRENGTH HUMBE	MIMBER
		•		•	1 = 46	8			
9.5 46.	6949.52	0 P. 1 2(40,5)	27339.521	56,566	3.2095-06	٦	6.673E+01	3	.
8.5 3	5868.07	P 1 2(39.	30.98	631:236	-342E-0	.975E-2	8 OE + O	.53296F-	~
.5 38	4796.10	P 1 2(3	9.9		145E-0	• 17 RE	3146+0	2328E-0	₽ 7 ;
6.5.3	3734.73	P 1 2(37.	97.4	583530	•959€-0	. 228E-2	677E+0	.51573E-0	(ئ ا خى
5.5 36.	2684.63		2	561-10	.230E-C	- 0 24E - 2	0725+0	.51036E-0	יי א
5 46.		P 2 7 (60.	8257.5	537.860	2 6	5.569F-22	2 + L.7	09727E	۰ ۲
0.5 40.	6918.10	0 1 2(40.	0.263.6	537.09	118E-0	.461E-2	.524E+0	-29096E-0	æ
•	0622.19	P 1 2(34	8410.84	.776	03E-0	07 9E-2	958E	5064RE-	σ,
8.5 39.	5837.35	(39.	8445.6	514.474	7265-0	- 082E-2	• 813E+6	.99824F+U	D .
9.5 39.	5837.35	39.	8451.69	513.723	•653E-0	. U36E-2	.629E+0	.26730E-0	= :
24.24 C.V.	5818.58 9611.74	0 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	28571.461	511.405	2 5	. 34 /E- 2	- COLUE - COLU	. 11" CCE) P
7.5 38.	24766-19	P 2 2(38.	8626.8	2.228	.056E-0	-916E-	.7052+0	89915E+0	<u>.</u>
8.5 36	24766.19	9 1 2(38.	8632.86	491.491	-234E-0	027E-2	. R22E+0	-2495FE-	15
8.5 39.	24747.28	P 1 1(39.	8651.	9-186	.455	• 539E-2	.867E+0	.01540E+3	16
1.5 32.	18616.42	P 1 2(32.	8732.00	9.443	909E-0	. 196E-2	. 043E+0	.51232E-0	17
6.5 37.	23705.52	P 2 2(37.	8801.50	71.646	.117E-0	. 237E-2	•627E+0	.80001E+0	
7.5 37.	23705.52	0 1 2037.	8867.51	n (- 505	.639E-2	• 114E • 0	0-310052	61
7.5 38.	23636.47	P 1 1(38.	A826.5	66.028	-334E-0	.715E-2	. 8 C 3E + 0	.91655E+U	92
0.5 31.	17637.22	15)2 1	8885.	77	2705-0	1.50 ZE - ZZ) (
7.7 Jb.	52.05022	P 2 2136.	8970'01 8975'07	50.166 50.166	-6396-E	. 197F-2		24281F-0	2 6
6. F. 37.	2002027	p 1 1(17.	8995.17	198	229F-0	531F-2	.743E+0	81766E+0	: 2
9.5 30	16675.09	P 1 2(30.	9038.4	2-729	.280E-0	. 682E-2	.382E+0	52875E-0	52
0.5 43	26949.5	0 2 2(40.	9120.4	ň	1-110E-03	- 00 BE-2	573E+0	14577E+	92
4.5 35.	21619.25	P 2 2(9132.6		2896-0	. 076E-2	- 563E+0	.60155E+0	. 12
5.5 35.	21619.25	0 1 2(35.	9138.5	9.89	.624E-0	.010E-2	047E+0	.23422E-0	8 2 (
5.5 36.	21599.89	P 1 1136.	9157.89	19	•160E-û	. 26 RE-2	.717E+0	.71876E • 0	62
.5 29.	15731	P 1 2(2	9186	3425-2545	.3195-0	3.96 CE-22	- C	1.54128E-31	0 F
	26.68602	2634	700000	, ,	5 7 7 C - 0	6.22F-3	676F40	2625550 2625550	100
4.7	C.CVCU7	- c	74-6676 9407-63		בי ענייי	2-15b	.336E+0	04633540	3.5
1.5 39.	25868.07	39	9313	5	.302E-0	95 DE-2	049E+0	.95112E-0	36
4.5 35.	20575.99	P 1 1(35	9314.	2	.806E-0	.132E-2	.721E+0	.61978E+0	
9.5 40.	25855-39	P 2 1140.	9323 •		-107E-0	. 93EE-2	879E+0	.52704E-0	9
0.5 40.	25855.39	0 1 1 (40.	9326.41	917	283E-0	. 925E-2	. 623E*U	* 1353/E*U	
7.5 25.		16.0212 1 7 0	29531.326	620	.941E-0	84.3F	.590E+0	.40283E+0	66
1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	19585.99	0 1 2 (33.	0.7446	952	.039E-0	.531E-2	.145E+0	.25739E-0	
3.5 34	19566.2	1 1(34.	9466	682	.733E-C	-713E-2	713E+	.520776+0	15
6.5 27.	13900 99	0 P 1 2(27.5)	9473.4	.904	.328E-0	·326E-2	•029E+0	.57587E-0	
8.5 38.	24796.18	0 2 2 (38.	9486.7	0.375	.869E-0	.052E-2	.159E+0	.74683E+0	
9.5 38.	24796.18	R 1 2(30.	9492.8	9.677	.399E-6	• 767E-2	7.26E + U	0-120/06.	
8,5 39.	24783.3	P 2 1139.	9499.	.	195E-0	2525-2	. 4 1 AF + 3	-367676. -93610F	ה על ה
9.5	24783.34	(3°C)	9565.7	202	.327F-6	25 CE - C	. 645E+0	333361.*0	
1.5 32.	20.19201	. 26136	0 1 0 1 0 1 0	4		ROOF-2	46.9F	280995	
	8591.E	n 1 2032 p. 1 2026	53354Y	76.5	-3206.	32F	141E+0	5983	*
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| TRANSITION | | 0 2 2(37. | R 1 2(37. | P 2 1(38. | p 2 2(31) | 0 1 2(31. | P 1 2(25. | P 1 1(32.
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| OMER ENERGY | 4 | 23734-737 - | 3734.73 | 3721.72 | 31613.10 | 7613.39 | 2154.87 | 7593.29
 | 2684.6 | 0.000 | 2671.5 | (652.2 | 6652.2 | 1315.6 | 1666.9 | 1546.0 | 1633.4 | 2.0072

 | 1633.4 | 0500.2
 | 5688.6 | 4785.26 | 0622.19 | 4785.26 | 9709.60 | 060 A.58 | 0608.58
 | 47.04.43 | 3881.3 | 5818.5 | 3351.32 | 9611.74 | 3860.20 | 9611.74 | 9597.90 | 2998 37 | 2998.3 | 8206.0
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| | JG LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY INTENSITY EINSTEIN A LINE STRENGTH HUMAN VAC CH-4 ATD ANGSTROMS CH-7 ATH-1 CH/MOLECILE SFG-1 | JO LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY INTENSITY EINSTEIN A LINE STRENGTH MUHNE VAC CM-1 VAC CM-1 T = 4600.0 K | JU JO LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY EINSTEIN A LINE STRENGTH HUHINF VAC CM-1 VAC CM-1 VAC CM-1 VAC CM-1 VAC CM-1 AIR ANGSTROMS CM-2 ATM-1 CM/MOLEGULE SEC-1 T = 4600.0 K T = 4600.0 K 7.5 37.5 23734.737 - 0 0 2 2137.5) 29650.278 3370.7716 8.5346-03 5.3496-21 8.041E+04 7.54726E+01 5 | JU JO LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY EINSTEIN A LINE STRENGTH HUMBE VAC CM-1 VAC CM-1 VAC CM-1 AIR ANGSTROMS T = 4600.0 K T = 46 | JU JO LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY INTENSITY EINSTEIN A LINE STRENGTH HUMIN VAC CM-1 VAC | JU JO LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY EINSTEIN A LINE STRENGTH HUNING SEC-1 VAC CM-1 VAC CM-1 VAC CM-1 VAC CM-1 AIR ANGSTROMS CH-2 ATM-1 CH/MOLECULE SEC-1 SEC-1 T = 4600.0 K T = | JU JU LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY INTENSITY EINSTEIN A LINE STRENGTH NUMBER OF THE CHYMOLECULE SEC-1 VAC CH-1 VAC CH-1 VAC CH-1 VAC CH-1 ANGSTROWS CH-2 ATH-1 CH-2 ATH-1 CH-2 ATH-1 CH-2 ATH-1 CH-2 ATH-1 CH-2 ATH-1 CH/MOLECULE SEC-1 T = 4600.0 K T = 4600.0 | JU JU LOWER ENERGY VAC CH-1 VAC CH-1 AIR ANGSTROMS CH-2 ATH-1 CH-2 ATH-1 | JU JU LOWER ENERGY VAC CH-1 VAC CH-1 T = 4600.0 K T = | JU JU LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY INTENSITY EINSTEIN A LINE STRENGTH HINTON VAC CM-1 ATR ANGSTROMS $CM-2$ ATM-1 $CM/MOLEGULE$ SEC-1 SC-1 SC-1 SC-1 SC-1 SC-1 SC-1 SC-1 S | JU JU LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY EINSTEIN A LINE STRENGTH HUMIN VAC CM-1 ATR ANGSTROMS CM-2 ATM-1 = CM/MOLEGULE SEC-1 | JU JU LOWER ENERGY TRANSITION FREQUENCY WAVELENGTH INTENSITY INTENSITY EINSTEIN A LINE STRENGTH HUMIN VAC CM-1 AIR ANGSTROHS CM-2 AIM-1 CM/MOLEGULE SEC-1 7.5 37.5 23734.737 - 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	LINE STRFNGTH NUMBER		.32200E-	.36194E-0	-11234		2286	389975-0	.50471E10	-310v0f-	.00093E • 0	.550684-0		BROZZE	. 53854F + 0	. 12877E+0	22001:	* * U * U U E - U	43203F-0	. 69967E+0	.76291F-0	. 16590F .	- 46498F-	77.	444075	84748C-	.7981BE+0	. 22180E+0	.00164E-U	200	43073E-0	.54887E-0	30	.02960E-0	.13746E+0	. 6964.2E+0	-210015-	72606E-0	.64211F-0	.82842	.12260E+0	118155-		.2067AE	
	EINSTEIN A SEC-1		.577E+0	.619E+0	1969B	753540	331E+	.166E .0	.405E+0	. 547E+	.515E+0	•	7		•	•	•		, ,,		٠,	3	•		•	•	•	8.361E+04	• •	•	: 7		7.392E+84		٦,	٦.	.116E+0	· CARETO	.704E+0	.821E+0	. 4595 + 0	.4.65E+C	3.695E+03	.439E+0	
	INTENSITY CH/MOLECULE	00.00 K	.586E-2	.873E-2	99.8E	• 645 55.7 Z	. 01 AE-	261E-2	.576E-2	. 322E-2	. 22 F.E-1	. 857E-2	3. 91 6E-19		. 81 9E-1	.26RE-1	. 48 EE - 2		515F-2	-495E-1	.123E-2	. 26 9E-1	•	. 6711E-2	1.357E-10	3.796E-2	1.1	2.233E-2	9	7. 000E-1	3.115E-2	4.7135-2	2.042E-20	758F-2	. 90 1E-1	.127E-1	. 826E-2	76 9E	. 62 FE-2	. 24 6E-1	33115-2	.0485-2	.101E-	052E-2	
	INTENSITY CH-2 ATH-1	9	.530E-0	.583E-0	-188E-	.5426-0	433E-	.012E-0	.368E-C	•	956E+0	.153E-U	•	576E-0	.093E-0	.022E	.371E-0	3-1956-04	799F-0		.174E-0	0-390 % ·	5.278E-03	.527E-U	i.	5.912E-06	5	rů.	٠,	년 -	• 5	₹.	3.258E-02	1875-0	.161E+0	.394E+0	.408E	3956-1	.057E	.584E+0	.267E-0	3-5482.	52E+0 28E-0	.869E-	•
•	WAVELENGTH IR ANGSTROMS		3.151	2.919	2.073	244.1	1.51	9.656	3.037	5.08	3.478	٠, ,	۰۰	, _	ė		٠.	, נ) LC			Ň	ᇵ.	i.	• -	• •			٠.			Ň	<u>.</u> .	: -	-			•			•	•	3181.6457		
	FREQUENCY VAC CM-1 A		31016.576	1018.80	026.95	27.0501		1050.25	1056,21	1004.68	1109.93	1114.07	31115.369	31.	113	38-83	165.00	185.58	191.89	202.18	206.15	214.47	39.6	231.10	231.64 276 25	36.	1291.	1293.	1295.	1636	1313.	1313.	9.5	3 4	1331.0	1377.	380.9	1 3 43	1000	1408.2	1413.	1 6		1439.	
	TRANSITION		R 1 2(27.	0 1 2(20.	P 2 1128.	K 2 2134.	P 1 1(2	n 2 1135.	R.1.1(35.	P 1 2(14.	P 2 2(19.	0 1 2(19.	11 2 2(26.	P 2 1(27.	0 1 1(27.	P 1 1(20.	R 2 2(33.	0 2 1134	p 1 2013.	P 2 2(16.	0 1 2(16.	0 2 2125.	R 1 2(25.	p 2 1(26.	p 1 1(19.	- N	P 2 2(17.	R 2 2 (32.	0 1 2(17.	p 1 2012.	0 2 1(33.	R 1 2(24.	R. 1 1(33.	0 p 2 (125,5)	0 1 1125	P 2 2(16.	0 1 2(16	0 2 2(23.	P 1 2121.	P 1 1(17)	R 2 2(31.	P 2 1124.	1 1124	7 -	9
	LOWER ENERGY VAC CM-1		6.006	8193.99	3885-2	7. T.	8169-86	0575.99	575.99	17.41	483.87	183.0	016-96	00.79	7-000	459.10	585.99	565 565	00000000000000000000000000000000000000	001.92	101.92	154.8	.87	138.24	67.76	7 · 60 60 60 60 60 60 60 60 60 60 60 60 60	6148.94	91.62	6	89.85	18571.728 -	315-66	21.175	122.52	298.5	525.69	25.6	5000.2	ב ע	24.664	513	04.82.5	82.5	7593	1
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<i>r</i> ,	LINE STRENGTH	.83093E- .59436E+ .57344E-	.54292E+0	.72807E.	.21720E-0	. 7357 FE . J . 0233 FE . O	1.49196F:01 7.9156F-01	0-Ju5664.	107335+0	.341371. .85058F-	.6275at +0	327886 -	103866-0	1.916611-01	302931-0	924075	0-1/0051	0+150925	545530-0	0.100700	33330E+0	32902E-	74207E-0	938256+0	2 1	82472F*0	589598-	-JU2665	80457F-	906201.	1061/1.	57791F-	736221 1	4.287815-01	74397F-	
	EINSTEIN A I	.296 .965	123E+9 082E+0	-942E+0	.853E+0	.930E+0	.067E+0	.588E+0	.530E+0	.351E * U .515E * O	-057E+0	3655+0	•1576. •695E•0		.716E+0	170E + 0	.072E*U	0	1.0375.0	063E+0 940E+0	384E+0	.352E+0	.005E+0	.786E+0	567E+0	.262E+0	. 604E + 0	. 112E+0	.394E+0	.177E * 0	. 366E • U		. 986E + 0	17E+0	.373E+0	
	INTENSITY CH/MOLECULE 00.0 K	0.100E-24 2.478E-10 1.142E-19	.134E	.641E-1	-112E-	.103E-1	60 -3	. 24 9E-2	48 5E-2	1.276E-20	.054E-1	. 523E-2	.339E-1	1.5316-23	. 564E-1	. 067E-2	74.7E	1-32.4.	. 08 4E-2	.474E-2 .063E-2	.917E-1	•	352E-1	.355E-1	2.359E-20	5-3199.	•	.601E-2	.679E-2	2126-2	54 JE	. 97 5E-1	. n74E-1	3.1645-20 6.2765-18	.691E-	
•	INTENSITY CH-2 ATH-1 T = 460	1.292E-05 3.953E+00 1.822E-01	.473	.214E+0	774E-0	.648E-0	4.527E+00	.157E-0	156E-0	.410E.0 .035E-0	.873E+0	.429E-0	.137E-6	43E-6	.974E-0	-096E-0	. 7 7 9 F - N	.540E+0	.729E-0	.033E-0 .291E-0	.059E+0	.603E-0	.752E-6	,758E+0	.764E-0	.541E-0	.422E-0	.554E-0	.274E-6	0-306-0	•12/E•0	.151E-0	.505E+0	5.047E-02	. NB BE-0	
(TABLE 9, cont. MAVELENGTH: AIR ANGSTROHS	78.328 77.690 77.335	175-31	174.479	173.469	170.95	169.	168.912	168.	166 - 117 y	166.332	165.279	15%.05B	162.4	161.597	160-206	159.012	150.503	150.232	157.669	157-100	154.653	154.235	152,303	151.876	150.347	150-269	149-043	147.74	147-72	95 · / 5 I	ج ج ا- ا	145.528	145.11	143-415	
	TAR FREQUENCY VAC CM-1 AI	1453,97 1460,39 1463,81	3.85	1492.11	1502.13	1527.11	31540.400	1547.44	1553	15	1573.13	1583	1595.73	1611	1620.	1633.54	1641.20	1651.34	1654.11	31659.760	1665.39	1690.	-	1713.65	31717.951	1733.34	1734.	1754-17	1759.59	17:19:73	1762.41	31765.021	1701.95	31786.009	1003.31	
	TRANSITION	R-2 1(38. P 2 2(15. Q 1 2(15.	22	P 1 1(16.	P 2 1123.	R 2 2130.	2, 2(14.	0 2 1(31	1 1(31.	R 1 26	p 1 1(15.	2 1122.	P 1 2(9.	R 2 1(37		. 2 2129.	N 2 2(20)	7 1 1(14.	n 2 1 (30.	1.630	0 1 1 (21.	1 2(8	0 1 2(12.	0 2 2(19.	2(19	R 2 2128.	P 2 1(20.	0 2 1129.	R 2 1136.	R 1 1129.	P 2 2(11	n 1 2(11).	0 2 2(10.	2(18.	P 2 1(19.	
	LOWER ENERGY VAC CM-1	1.72 2.93 2.93	709.	53.14	691.32	652.	9 9	631.94	6631.94	964.5	341.84	62	965.56 056.56	671.4	19-14	709-23	206-02	810.75	688.67	15680.670 +	186.35	16.169	44.52	48.464	10.161	A5.26	1,74.37	747793	633.41	4764.43	880.52	75.080	111.85	= :	790 .49	•
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LINE STRENGTII I	.91665E+0	.50864E+D	51668E+0	40599E-0	43172F+0	172685+0	.38879E+0	2.3271CE+01	71268E+0	78268F-0	.3117CE+0	.06410E-0	40874E+0	.33399E+0	4 8361E+0	526435-0	135546-0	50815E+0	.10604E+0	.46105E-0	23478E+0	.55838E+0	.93759E-0	. 8 9 9 6 2 E + 1	9 0938E-0	13788E-0	30850E+0	44981E+0	10162C+0	, 57735E+0	.69238E+0	41015E-0	.12757E+9	78036E-0	0+100760		27171E-0	.10199E+0	.48508E+D	.1218/E+D	.20814E+0
EINSTEIN A SEC-1	.739E+0	.631E+0	• 0 0 2E + 0	.175E+0	. 236E + 0	.997E+0	.991E+0	1.8526+05	3+U888	.311E+0	.120E+0	.447E+0	.131E+0 .746E+0	.819E+0	.037E+0	737F+0	.357E+0	. 025E+0	.220E+0	.711E+0 .961F+0	.835E+0	.377E+0	.063E+0	.139F+1	0416+0	•563E+0	453E+0	.848E+0	.114E+0	.875E+0	.3426+0	469E+10	.068E+0	.063E+0	. 235E + 0	0000	.160E+0	.970E+0	.667E+0	*555E+0	• 959E + 0
INTENSITY CH/MOLECULE 00.0 K	.528E-1	.241E-1	. 050E-1	. 395E-1	. 514E-1	.074E-1	.261E-1	4.354E-19	28 UF-1	1195-2	.435E-1	. 905E-1	. 291E-1	. 148E-1	1069E-1	127F-1	3485-1	. 992E-1	. 90 BE-1	. 373E-1 . 641F-1	. 50 1E-1	.197E-1	. 642E-1	587F-1	9316-1	.320E-2	585E-1	. 744E-1	.605E-1	. 27 4E-1	.329E-1	253E-1	.2152-1	.723E-2	- 039E-1	7 - 36 D 2 - 12 - 12 - 12 - 12 - 12 - 12 - 12	. 811E-2	.264E-1	. 24.5E-1	• U41E-1	• 174E - 1
INTENSITY CH-2 ATH-1 T = 46	-042E+	.170E	•093E	-417E-	. A 32F	.095E+	.478E+0	94	1635+0	.786E-0	.186E+0	.039E-0	.346E-0	.022E+0	.705E+6	393F-0	.341E-0	.275E+0	.262E+0	.786E-0	.99.0E+6	.910E+0	.215E-G	.371E+6	.676E-0	.107E-0	.910E	.782E+6	.156E-0	•033E+0	.329E+C	1916-0	.151E+0	.748E-0	0 * 1255°	023010	.689E-0	.478E+0	.986E+0	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	•145E+0
WAWELENGTH Air Angstroms	4.76	. 61	3.37	3, 03	7	2.09	2,36	3110.5270	27	9.56	3.34	9.6	900	7.56	7.46	١, ٨	5.05	5.01	5.67	2 to 20	3.35	3.27	2,39		13	2.13	1.66	9.60	3.58	3.54	3.42	3.81	3.72	9.69		. u	9.0	6,37	5.33	5.12	0.0
FREQUENCY VAC CH-1	2095	2097	2110.20	2113,21	2122.40	2123.40	2123.68		2142.74	2149.54	2151.82	2154.64	2167.15	2170.23	2171	2483	2185	2186.	2189	2192	2213.8	2214.	2223	50.4222	2226.49	2226.52	2231	2252.86	2253.	2253.	2254	2261.	2262.0	2562.2	2265.	22010	2285.55	2286.	2286	6822	523
TRANSITION	1 1616	R 1 1(2	0 2 2(12.	P 1 2(3.	R 1 2(12.	P 1 11 7.	0 1 2(5.		0 1 1(13.	R 2 1(33.	0 2 2(11.	R 1 2(11.	R 1 1124	P 2 21 4.	0 1 20 4.	P 2 1612.	P 1 2(2.	0 1 1(12.	n 2 2(10.	7 1 2(10.	P 2 21 3.	0 1 24 3.	P 2 1(11.	0 2 2 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4	R 1 2(9.	0 2 1123.	R 1 1123.	P 2 2(2.	P 1 2(1.	0 1 21 2.	0 2 2 6 8.	P 2 1(10.	R 2 2121.	2 1(32:	0 1 1(10.	1 2 2 1 0	2 1122	P 2 2(1.	1 2(1.	1 16	1 1622.
ONER ENERGY VAC CH-1	4.156	277.107 -	8.880 +	429.275 -	4 4 20 4	26.730 -	824.525 +	484.962 +	319,355 4	602,331 +	2884.110 -	2884.110 -	62,394 +	608-188 -	08.188	57,435	288.769 +	852.465 -	4 641.254	53.149 +	459,459	459.458 +	419.051 +	170.000 -	56.568	672.362 -	672,362	19.041	67.491 -	289.041 -	94.914.4	019.633 -	931.446 +	622.844 -	114,655 ~	350.720 -	+ 868.406	187.752 +	7.751	355.10	69.7
96 11	3	•	ć	,	•	• •	•		,	m	-	.		3	 	• •	;	2	Ġ	•	رم د	1	•	•	• 6	٠	m 4		•	•	•			è.		• 1	• •	7	•	•	•
. 3	3	26.95	è	•	•	• •		24.5	, v		1.	å.	2 R	100	•	• :		2	ċ	11.5	, ,	1	10.5	5 +		m	÷.	• •		•	60 0		2		.,	•	• •	•	1.5	٠	23.5

UHBER.	351	352	353	354	355	356	- E	M (5)	360	361	362	202	304	366	367	368	369	370	571	275	475	375	376	377	378	379	380	382	383	482	385	000	388	300	396	165.	392	56.5	e la	β. 6 0 M	397	Š	399	0	
LINE STPENGTH HUMBER	.95546E-0	89017E+0	.27554E+0	5922E-0	3333E+D	02768E+0 06667F+0	2470RF-0	58667E-0	58235E+0	58148E+1	42158E-		1324061	3333E+0	66645E-0	52912E+0	02998E+0	31818E-0	752536+	3333UET	31633E1	47348E+	92771E+	78065E-	01636E+	26352E+	58923E-	0 1696E+	11311E¢	U 5266E+	82765E+	2124	41497E+	776735-	90611E+	35645E+	.33651E+	3144CE+		n c	25111F	-78368E-	- 96646E-	6467E+	
EINSTEIN A L Sec-1	•999E+	.304E+	• 332E+	• 682E •	•810E+	167E+ 262E+	SAGE	702E+	.347E+	*095E+		• / / / E •	* 1000 F *	8265+	. 2045+	.843E+	.315E+	656E+	•974E+	1000.	.22164	6.357E+05	. 265E+	.894E+	• 995E•	. 317E+	.620E+ .777F+	. 160E+	.960E+	.216E	.357E+	1.2162E+05	• 023E+	+3605.	. 254E+	.611E+	•618E+	• / Car •		711E •	159F+	797E+	.521E+	.735E+	•
INTENSITY CH/HOLECULE 10.0 K	3.9796-19	293E-	8195-1	196E-1	142E-1	20	12001	132E-1	31.96-1	215E-1	63E-	06 9E - 1	44 CE - 1	1445-1	707E	038E-1	883E-1	822E-1	312E-	7 - 27 60	2016-1	9.083E-18	127E-1	. 613E-	. 95 8E-1	. 553E-	31 6E-2 76.8E-1	133E-1	. 247E-1	.728E-1	37 9E-1		. 612E-1	.4236-2	3942-1	•111E-1	117E-1	. 24 bt - 1	- 2299.	1.156E-18	685F-1	. 90 3E-2	36	-403E-	
INTENSITY CH-2 ATH-1 T = 460	•	483E+0	.248E+	.291E-0	.823E+0	448E+0	777F-0	709F-0	.487E+0	.916E+0		•128E+U	.491E+0	.825F+0	.105E-0	.038E	.253E+0	.289E-0	689	9641646	043162°	• •	.799E+B	.169E-C	.110E+0	.365E+0	5.291E-62	.808E+6	.316E+G	*233E+0	.201E+0	# # # #	.055E+0	.057E-6	.224E+0	.368E+	.782E+0	.369E+6	* 6566*	200	SAGF+0	.227E-0	.311E	.238E+0	
MAVELENGTH IR ANGSTROHS	0.95	095.341	194.626	094.460	127.860	193,61	75.50	192.58	192.39	091.36	1.3	091.24	160 ac		90 . 45	90:37	90.27	89.900	6.0	89.85 60.85	10.60	0.00	99.01	83,58	87.48	87.34	3086.6491	86. 22	35.32	85.1	84.89	3 M	83.	82.4	82.06	81.667	81.625	10	61.62	6 4	79.95	79.50	73	78.47	
FREQUENCY VAC CH-1 AI	2295.05	2297	2304.69	2306.4	2314.0	2315.2	0.363.0 5.363.0	0.5020	2328.0	2338.7	233	2340.0	2343.5	744.0	2348.2	2349.1	2350.2	2354.1	2354.4	2.426.5	7.354.7	235	2363.4	2366.82	2379.40	2380.91	238	2392.64	2402.12	2403.40	2406.57		2423.57	2432,07	2436.36	2440.54	2440.98	2441.80	99.4442	32455.599	2454 556 2458 566	2463.35	2471.1	2474.17	* ******
TRANSITION	. 6 11 6	1 1 0	0 2 21 6.	R 1 21 6.	0 1 26	R 2 2(20		2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	0 1 1 0	0 2 2 6 4	2 1(21.	R 1 2(4.	P 1 1 3	7 1 11C1.	8 1 2 E	0 2 21 3.	R 1 2(3.	P 2 11 7.	0 2 2 1	0 2 21 2.	K 1 2 1 1	K 1 C1 C	R 2 2(19.	R 2 1(31.	P 2 11 6.	0 1 16 6.	0 2 1(20	R 1 1(20)	P 2 1(5.	0 1 16 5.	R 2 2(18.	R R 2 2(.5)	0 1 16 4	0 2 1(19,	R 1 1(19.	P 1 1(1.	P 2 1 (3.	0 1 16 3.	R 2 2(1/c	R 2 21 1	0 1 11 2	R. 2 1136.	0 2 1(18	P 2 11 1.	
LOWER ENERGY	1654.577 +	5	78.	1076.509 +	26.44	93.9	19. 42	10.479	1324.201	608.19	69.86	08.19	101.92		126.291 +	9.27	29.27	29.0	65.40	88.76	64.18	+ 650-620+	483.87	60.41	694	69.21	7459.105 +	11.00	00.3		01.92	126.449 -	55.9	76.	776.4	0.00	-	202.37	76.0	<u> </u>	200	716.07	6122,62	.05	
- Pr	7. O. 7.	1 . C . C . C . C . C . C . C . C . C .	5.5 6.5	_	.5 .5	2	-	٠	מים מים	, W	.5	7 5		ئ در			e.		S.	ر ا	r.	7.5 2.5		, ול הו	2	rč.	5 20			, s	.5 18.	n n		5 19.	0.5 1	•	.5 3.	ر س	.5	4		1.5. 76.	18.5 18.5	.5	

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	ភូមិពល	401	20%	£03	2 L		107	£0.3	601	r. 10		212			416	417	ю : :	61.	4,24	221	121	121	521	26	128	1,29	430	11.5	1 2 2	4, 34	135	277	1,38	613	D • 3	245	£443	3 2 2	5	E .	· ·	C (**)	450		
	LINE STRENGTH HUNGER	2.24821E+60	0 2 0 8	723(+	י כ	1-112122	1.526art 193	206545+0	42637F+0	- 39 20 8	02.3	0 1 1 1 1 1 1 1 1 0 0 0 0 0 0 0 0 0 0 0	76.97	25771	0	5.178210 100	1.273056 + 30	9	1.500625+01	ころ	184	1.34381F . 00	2.525051.400	I. 12416F 101	7.11106E-01	1.02315E+01	1.398600+01	0.20415[+00	1.317336 00	3.559400 00	7.485415-01	1.251096106	613	7.90527E-01	172	71875	8.38010E-01	093 -E+0	9647E+0	72852E • 0	-91565E-	9.51813E-01	86803E+0		
	EINSTEIN A SEC-1	4+108E+05	• 342E+	523E+0	. U58E * U	- / U / E • U	2.424E+05	.313E+0	. 658E+1	.109E+0	.498E+0	2.498E+05	0.00000	711E .0	. 803E+0	.629E+0	57E+0	2715.0	564540	1565+0	.717E+0	228E+0	. 831E+0	.791E+0	. 464E + 0	610E+0	619E+0	806E+0	577F + 0	137E+0	0+3569	7.5275+04	352E+0	975E+0	9896+0	2.502E+05	317E+0	835E+0	71BE	504E+0	0+32+2	3.953E+04	722E+0	•	
	INTERSITY CM/HOLECULE GDO.O K	2.010E-18	•	. 972E-1	635E-1	5405-2	2.013E-18	467E-1	641E-1	9.912E-20	36.3E-1	3.019E-13	1 - 2 - 6 - 6 - 6 - 6 - 6 - 6 - 6 - 6 - 6	.97.E-1	-764E-2	. 46 OE-1	1.107E-18	37.75	72/12-1	297E	.772E-1	7	1-3060	. 582E-1	3.955E-18 1.626E-19	. 807E-1	3.091E-18	. 014E-1	. 920E-1	7865-1	1-35,00	3.441E-18 9.24FF-19	.371E-1	. 557E	. 997E-1	3. A17F-18	15.05		. 026E-1	195-1	.892E-1	5.7036-19	21.35-1	,	
	INTENSITY CH-2 AIM-1 T = 4	.207E	694E+0	.146E+0	0.4056.	2195-1	3.212E+00	936E+G	214E+6	1.5A1E-01	•	4.817E+60	0 1765-01	4.751E+00	9-197E-04	5.521E+00	1.765E+00	00124242	346546	5.261E+00	.019E+C	.789E+0	.334E+0	.714E+6	5.505E-01			-	•	4.455+00	-	1.475F+00	.378E+0	-080E-0	.276E+0	5.9995+80	.055E-0	.085E+1	.424E+0	•	209E-	0-36-0	22F+)	
in 7) cours.	WAVELENGTH AIR ANGSTROHS	•	_		3677.0346	Ž (3075.1188	376	031	072.66	•	207		071.148	071.	070.	395	2 6	, 5	3069.6780	3069-1840	3068.0053	3068.7061	3068-6120	3068-2359	3067.9436	3067-9124	8	300/-00/0	3067.2418	3066.6138	3066.3109	3065.9766	3065.3718	3065.1128	3065,0879	69.990	376	064.230	064-190	063.971	3063,9256	63.726		
	REQUENCY VAC CM-1	474	2475.3	4.8742	2489.3	2507.2	32569.624	2517.4	2531.	25	2539.4	1.0 to	1467	2551.6	2552.13	2558.6	2559.63	2260	2201062	2567.24	4.2725	2576.50	2511.55	2578.5	32582.019	2585.65	2505.98	2587.37	258845	01.6652	2599.78	2603	2606.55	2612.99	\$615.74	2616.0	2622.32	2523.60	2625.13	6-529	2627.90	628	7630 64	70.000	
	TRANSITION	0 0 1 1(1.5)	R 1 1(16.	R 2 2(16.	R 2 21 2.	2111 2 0	R 4 1117.	R 2 2(3	R 2 2(14.	0 2 1116.	œ		7 7 7 7 6	R 2 2(13.	R 2 1(29.	R 2 2(5.	0 2 11 2.	K 1 11 C	C111 7 0	R 2 2112.	R 2 21 6.	0 2 16 3.	R 1 11 3.	R 2 2(11.	<u>د</u> د	R 2 2(10.	R 1 1(14.	R 2 2(8.	R 2 21 9.	2 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7	0 2 1(13.	7 1 1(13	R 1 11 5.	0 2 1(12.	9 11 2 1	1612		0 2 11 7.	R 1 1(11.	œ	C	0 2 11 8.	5 11 2 2	• ATL T V	
	LONER ENERGY VAG CH-1		122.6	25.69	50.682	498 645	5498-455 +	429.45	1.36	904.62		608-188 -	9 6	0.00	790.64	24.52	83.719 +	17.50	10 · 11	7.5	077.85	201.92	201.92	580.52	1367-617 +	450.27	97.019	3.29	054.35	5.10	1.98	317	75.5	6.0A	767.45	6.03	4.107	10.01	3.61	026.73	015.03	N 1	650 • 7 q	+ 950.51U2	
•	JO 76 1	1.5 1.5	.5 16.	.5. T	2,	֓֞֜֜֞֜֜֞֜֝֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓	18.5 17.5	5.	.5 14	.5 16.	.5 16.				.5 29.	.5 5.	2.5 2.5		.7 15.	, ,	5.6	m	.5 3.	.5 11.		5 10.	-	.5 8.	ب و.	7.7. c.3.	.5 13.	.5 ±33	, r,	.5 12.	.5 6.	12.	2	- -	. 5	•	10.5 16.5	8.5 8.5	v.	5.01 5.11	

				TA	ABLE 9, cont.					• >>>
3	JG LOWER	LOWER ENERGY Vac CH-1	TRANSITION	FREQUENCY VAC CH-1	MAVELENGTH AIR ANGSTROMS	INTENSITY CK-2 ATH+1 T = 40	INTENSITY CHZMOLEGULE 4600.0'K	ETHSTEIN A I SEC-1	LINE STPFNGTIF NUMNFP	mnfr ,
9.5 €	5.	321.252 +	R R 1 1(8.5)	32630.565	3063.7208	6.829E+00	4.28GE-18	2.671E+05	7.7 40501 + 00	151
٠.	.5	1650.790 -		32632.269	30.63.560.0	•	4.305E-18	2.707E+05	8.827CCT 1 00	1.00
.5.2	~	5.2	R 2 1 (2	32533,357	3063.4586	.337	8.382E-22		1.79A TAF-01	1,77,4
ı,	.5	*020*	R 2 11 1.	2643.	3062.5260	2.672E-01	1.674E-13	.305E+0	1.87037F-01	25.2
ħ.	•	3	R 2 16	94.50	3057.7289	•	2.592E-19	•	2.98634F-01	200
.5 27	.5 13	000	R 2 1127.	2707.	3056.5390	1.927E-03	1.20 RE-21	.008E+	1.01036F-01	456
2	•	202.370 +	R 2 1(2744.5	3053.0544	4.629E-01	2,901E-19	•	3.488785-01	159
5.2	• 5	-	0.	73.9	3050-3154	2.745E-03	1.721E-21	1.140E+03	1.82567E-01	16.54
Ŗ,	L.	55.	R 2 11 4.	26 22	3049.57.04	4.564E-01	2. 86 1E-19	.221E+	53	459
•	.5 11	2	₽′	32033.743		•	2.421E-21	•	844518-	4.50
r.	.5	, i.	R 2 1(32838.395	3044 - 3301	4.219E-01	2.644E-19	1.873E+04	3.594965-01	192
•	ī,		2	320A1.122.	3040.3740	3.756E-01	2.354E-19	.569E	3.47219F-01	29 %
.5 2	5 1	482.58	R. 2. 1(24.	2886.	3039.8524	•		* 4,6E	1.867126-01	1,63
.5	r.	0.620	R 2 11 7.	6262	3036.7331	3.267E-01	2.047E-19	1.3195.04	3.316531-01	1,51,
.5 23		591.3	R 2 1(23	2933.	3035.5650		4.654E-21	•	1.893028-01	ני גינו
=	٠. د	124.29	R 2 1(0.	2956	3033.4317		1.751E-19	1.1155.0%	3.15383E-01	466
.5.25	.7.	926	R 2 1122.	2973.	3031-0775		6.350E-21	1.019E+03	1.924711-01	1273
.5	.5	654.5	~	2903.	3030-4901		1.47 KE-19	9.494E + 0 3	2.996535-01	4.6.8
21	.5	1 n 6.	R 2 1121.	3.007	3028.7695			2.041E+03	1.96077F-01	t () t
.5 16	r.	10	R 2 111	33616.299	3027 - 9254	•	1.230E-19	8. 243E+03	2. N5026F-01	0.4.0
.5 20	_	474.37	R 2 112	3034=08	3926.2215		1-14RE-20	2.286E.03	2-001646-01	471
-5 11	5	19.08	R 2 1(11.	3639.	3025.7534	1.618E-01	1.014E-19	7.035E+03	2.716731-01	1,72
13	ıv.	6 1 0	R 2 10	3056.86	3024-2155	2.426E-02	-52CE-	2.565E÷03	16-159050-2	673
.5: 12	ı,	55.	R 2 1(12.	3059.2	3023.9887	1.320E-61	8.274E-20	∹	2.596266-01	6,7%
19.5 18	אָן	2	~	3073.	3022.7341	3.101E-02	1.994E-20	2.011E+01	2-10176F-01	513
	• 5	3319.355 +	R 2 1(13.	3073.9	3022,6451	.057E-0	6.686E-20	5.338E+03	2.40029E-01	476
.5.1	S	10.15	~	330.83.652	3021.7608	4.132E-02	2.590E-20 ·	3.241E+03	2.161A3F-01	477
15.5 14	5.	3819-156 -	R 2 1(14.	3083.9	•	-	5.353E-20	4.687E+03	2.39187F-01	478
.5.	•	- 161-516+	R 2 1(16.	3066.91	3021.2801	5.316E-02	3.332E-20	3.655E+03	2.22962E-01	479
•		4351,310 +	R 2 1(15.	3008.9	3021-2768	6.766E-02	4.241E-20	4.129E+03	2.30598E-01.	4 8 0

THE INTEGRATED INTENSITY FOR THE DAND IS THE EINSTEIN A CORF FOR THE BAND IS

8.68628E+02 CH-2 ATM-1 AT T 7.07924E+07 SEC-1

5.44436E-16 CM/MOLECULE

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4600.0 K

ON POOR PROKES

California de la calactería de la calact

TABLE 10

T (K)		c _R		T (K)		Q_{R}
200	"X _a	26.71		2000		298,23
240		32.24		2500		375.87
296		40.17		3000		454.19
300		40.75	•	3500		533.18
500 ***	•	70.00		4000	V-	612.85
750		107.35		4500		693.14
1000		145.09	•	4600		709.26
1250		183.08		5000		773.98
1500		221.28		5500		855.28
1,750		259.67		6000		936.88